

Maxwell's equations

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Maxwell

(13 June 1831 – 5 November 1879) was a Scottish physicist.



Famous equations published in 1861

Maxwell's equations: integral form

$$\oiint \mathbf{E} \cdot d\mathbf{S} = \frac{Q(V)}{\epsilon_0} \quad \text{Gauss's law}$$

$$\oiint \mathbf{B} \cdot d\mathbf{S} = 0 \quad \text{Gauss's law for magnetism: no magnetic monopole!}$$

$$\oint_{\partial\Sigma} \mathbf{B} \cdot d\boldsymbol{\ell} = \mu_0 I + \mu_0 \epsilon_0 \iint_{\Sigma} \frac{\partial \mathbf{E}}{\partial t} \cdot d\mathbf{S} \quad \text{Ampère's law (with Maxwell's addition)}$$

$$\oint_{\partial\Sigma} \mathbf{E} \cdot d\boldsymbol{\ell} = - \iint_{\Sigma} \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S} \quad \text{Faraday's law of induction (Maxwell–Faraday equation)}$$

Relation of the speed of light and electric and magnetic vacuum constants

$$c = 1/\sqrt{\epsilon_0\mu_0} = 299792458 \text{ m/s}$$

$$\epsilon_0 \approx 8.854 \cdot 10^{-12} \text{ As/Vm,}$$

$$\mu_0 = 4\pi \cdot 10^{-7} \text{ Vs/Am,}$$

ϵ_0 permittivity of free space, also As/Vm or F/m (farad per
called the electric constant meter)

μ_0 permeability of free space, also Vs/Am or H/m (henry per
called the magnetic constant meter)

Differential operators

$\nabla \cdot$ the divergence operator $\nabla \cdot \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$
div

$\nabla \times$ the curl operator $\nabla \times \vec{A} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A_x & A_y & A_z \end{vmatrix}$ Other notation used
curl, rot $\partial_x = \frac{\partial}{\partial x}$

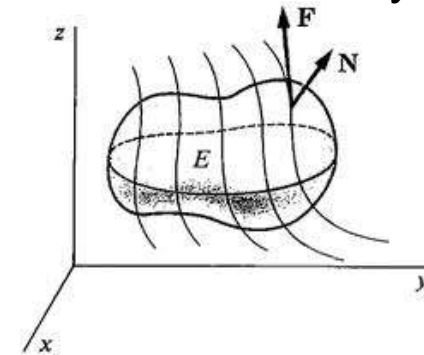
$\frac{\partial}{\partial t}$ the partial derivative with respect to time

Transition from integral to differential form

Gauss' theorem for a vector field $\mathbf{F}(\mathbf{r})$

$$\iiint_V (\nabla \cdot \mathbf{F}) dV = \oiint_S (\mathbf{F} \cdot \mathbf{n}) dS.$$

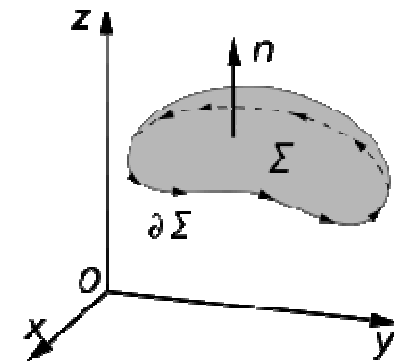
Volume V , surrounded by surface S



Stokes' theorem for a vector field $\mathbf{F}(\mathbf{r})$

$$\iint_{\Sigma} \nabla \times \mathbf{F} \cdot d\Sigma = \oint_{\partial\Sigma} \mathbf{F} \cdot d\mathbf{r},$$

Surface Σ , surrounded by contour $\partial\Sigma$



Maxwell's equations: integral form

$$\oiint \mathbf{E} \cdot d\mathbf{S} = \frac{Q(V)}{\epsilon_0}$$

Gauss's law

$$\oiint \mathbf{B} \cdot d\mathbf{S} = 0$$

**Gauss's law for magnetism:
no magnetic monopoles!**

$$\oint_{\partial\Sigma} \mathbf{B} \cdot d\boldsymbol{\ell} = \mu_0 I + \mu_0 \epsilon_0 \iint_{\Sigma} \frac{\partial \mathbf{E}}{\partial t} \cdot d\mathbf{S}$$

**Ampère's law
(with Maxwell's addition)**

$$\oint_{\partial\Sigma} \mathbf{E} \cdot d\boldsymbol{\ell} = - \iint_{\Sigma} \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{S}$$

**Maxwell–Faraday equation
(Faraday's law of induction)**

Maxwell's equations (SI units) differential form

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0} \quad \rho \text{ density of charges}$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$

$$\nabla \times \mathbf{B} - \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \mu_0 \mathbf{j}$$

\mathbf{j} density of current

$$\nabla \times \mathbf{B} - \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} = \frac{1}{c^2 \epsilon_0} \mathbf{j}$$

$$c^2 = \frac{1}{\mu_0 \epsilon_0}$$

Electric and magnetic fields and units

E	electric field,	<u>volt per meter</u> , V/m
B	the magnetic field or magnetic induction	<u>tesla</u> , T
D	<u>electric displacement</u> <u>field</u>	<u>coulombs per square meter</u> , C/m ²
H	magnetic field	<u>ampere per meter</u> , A/m

Constitutive relations

These equations specify the response of bound charge and current to the applied fields and are called constitutive relations.

$$\mathbf{D}(\mathbf{r}, t) = \varepsilon_0 \mathbf{E}(\mathbf{r}, t) + \mathbf{P}(\mathbf{r}, t)$$

\mathbf{P} is the polarization field,

$$\mathbf{H}(\mathbf{r}, t) = \frac{1}{\mu_0} \mathbf{B}(\mathbf{r}, t) - \mathbf{M}(\mathbf{r}, t),$$

\mathbf{M} is the magnetization field, then

$$\mathbf{D} = \varepsilon \mathbf{E}, \quad \mathbf{H} = \mathbf{B} / \mu$$

where ε is the permittivity and μ the permeability of the material.

Wave equation

$$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$

$$\nabla \times \mathbf{B} - \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} = \frac{0}{c^2 \epsilon_0} \mathbf{j}$$

$$\nabla \times (\nabla \times \vec{B}) = \nabla(\nabla \cdot \vec{B}) - \nabla^2 \vec{B} = -\nabla^2 \vec{B}$$

$$\nabla \times \frac{\partial}{\partial t}(\vec{E}) = \frac{\partial}{\partial t}(\nabla \times \vec{E}) = \frac{\partial}{\partial t}\left(-\frac{1}{c^2} \frac{\partial}{\partial t} \vec{B}\right) = -\frac{1}{c^2} \frac{\partial^2 \vec{B}}{\partial t^2}$$

Double vector product rule is used
 $\mathbf{a} \times \mathbf{b} \times \mathbf{c} = (\mathbf{a} \cdot \mathbf{c}) \mathbf{b} - (\mathbf{a} \cdot \mathbf{b}) \mathbf{c}$

$$\nabla^2 \vec{B} - \frac{1}{c^2} \frac{\partial^2 \vec{B}}{\partial t^2} = 0$$

$$\nabla^2 \vec{E} - \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

2 more differential operators

$$\nabla^2 \text{ or } \Delta \text{ Laplace operator or Laplacian } \nabla^2 \vec{A} = \frac{\partial^2 A_x}{\partial x^2} + \frac{\partial^2 A_y}{\partial y^2} + \frac{\partial^2 A_z}{\partial z^2}$$

d'Alembert operator or d'Alembertian

$$\square = \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial x^2} - \frac{\partial^2}{\partial y^2} - \frac{\partial^2}{\partial z^2} = \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2 = \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta.$$

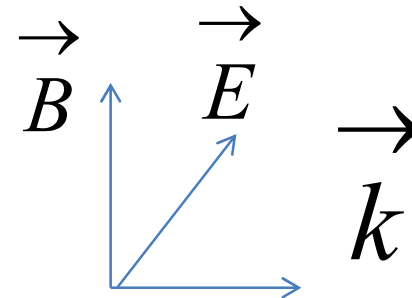
Plane waves

Thus, we seek the solutions of the form:

$$\vec{B} = \vec{B}_0 \text{Exp}[i(\vec{k} \cdot \vec{r} - \omega t)]$$

$$\vec{E} = \vec{E}_0 \text{Exp}[i(\vec{k} \cdot \vec{r} - \omega t)]$$

From Maxwell's equations one can see that



$$\nabla \times \vec{B} = i \vec{k} \times \vec{B} \quad \text{is parallel to } \vec{E}$$

$$\nabla \times \vec{E} = i \vec{k} \times \vec{E} \quad \text{is parallel to } \vec{B}$$

Energy transfer and Poynting vector

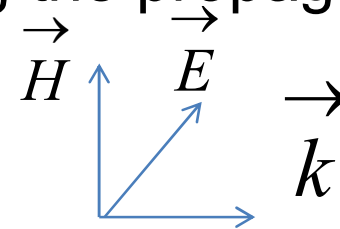
Differential form of Poynting theorem

$$-\frac{\partial u}{\partial t} = \nabla \cdot \mathbf{S} + \mathbf{J} \cdot \mathbf{E},$$

$$u = \frac{1}{2} (\mathbf{E} \cdot \mathbf{D} + \mathbf{B} \cdot \mathbf{H}). \quad u = \frac{1}{2} \left(\epsilon_0 \mathbf{E}^2 + \frac{1}{\mu_0} \mathbf{B}^2 \right)$$

u is the density of electromagnetic energy of the field

$$\mathbf{S} = \mathbf{E} \times \mathbf{H} \quad \parallel \quad \mathbf{n} = \frac{\mathbf{k}}{k} \quad \text{S is directed along the propagation direction}$$



Integral form of Poynting theorem

$$-\frac{\partial}{\partial t} \int_V u dV = \oiint \mathbf{S} \cdot d\mathbf{A} + \int_V \mathbf{J} \cdot \mathbf{E} dV$$

Energy quantities continued

$$\vec{B} = B_{\max} \exp[i(\vec{k} \cdot \vec{r} - \omega t)]$$

$$\vec{E} = E_{\max} \exp[i(\vec{k} \cdot \vec{r} - \omega t)]$$

$$B_{\max} = E_{\max} / c$$

Observable are real values:

$$\vec{B} = B_{\max} \cos[i(\vec{k} \cdot \vec{r} - \omega t + \Delta\varphi)]$$

$$\vec{E} = E_{\max} \cos[i(\vec{k} \cdot \vec{r} - \omega t)]$$

$$I = S_{av} = E_{\max}^2 / 2c\mu_0 = cB_{\max}^2 / 2\mu_0 = cu_{av}$$

$$u_{av} = E_{\max}^2 / (4c^2\mu_0) + B_{\max}^2 / 4\mu_0$$